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Parallel Natural Convection Simulation on GPU Using Lattice Boltzmann Method

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ABSTRACT: In this paper, we propose the implementation of a parallel lattice Boltzmann method (LBM) algorithm for the simulation of two-dimensional natural convection heat transfer problems. The LBM code is written using NVIDIA C language and is run on the graphical processing unit (GPU) of a Nvidia GeForce9600 card. The behaviour of the convective flows, which are driven by buoyancy forces, was numerically studied for Rayleigh numbers (Ra) from 1×10^3 to 1×10^6 . Our numerical results show a good agreement with the experimental and numerical results obtained in the literature. Based on performance comparison, it is shown that performance utilizing the processing power of the GPU is significantly faster than unaided CPU processing.

1. Introduction

Natural convective flows occur in many industrial situations. These kind of flows are particularly important in passive cooling electronic components, heat exchanger, solar energy and thermal storage systems. Such flows are very complicated. Therefore, it is necessary to clarify the flow structure and the mechanism in detail to establish an improved design technique. The numerical method has a powerful ability to investigate the flows and to clarify the predominant factors that influenced the flows.

Extensive numerical research works have been done in this field. Most of them are based on traditional numerical methods such as the finite difference, finite volume and finite element methods for solving the Navier-Stokes equations. De Vahl Davis (de Vahl Davis, 1983) provided "benchmark solutions" that were obtained with a second-order finite-difference scheme. The Navier Stokes equations are nonlinear equations and hard to be solved by traditional methods numerically, so recently the Lattice Boltzmann method (LBM) was introduced as an alternative method for solving Navier Stokes Equations especially in natural convection (Dixit and Babu, 2006; Guo, Shi, and Zheng, 2002). Lattice Boltzmann algorithms are simpler and easier to be implemented than the traditional methods. The main advantage LBM is its data parallel nature, so one can parallelize the LBM code relatively easily and accelerate the execution time.

In this paper, we propose the implementation of a parallel LBM scheme for the simulation of two-dimensional natural convection heat transfer problems on a graphical processing unit (GPU) (Obrecht et al., 2011). The LBM code is written using NVIDIA C language and is run on the graphical processing unit (GPU) of a Nvidia GeForce9600 card.

2. Problem Description

The physical domain of problem is described by figure 1. The domain consists of a square enclosure with side length of 1. Constant uniform temperatures T_h and T_c ($T_h > T_c$) are imposed at the left and right walls of the domain respectively. Buoyancy forces drive the circulation flow in the enclosure. These forces are induced by constant heating from the left side. An appropriate scaling for the flow regime is characterized by the Prandtl (Pr) and Rayleigh (Ra) number. In LBM, both of the parameters are related to density relaxation τ_ρ and temperature relaxation τ_g .

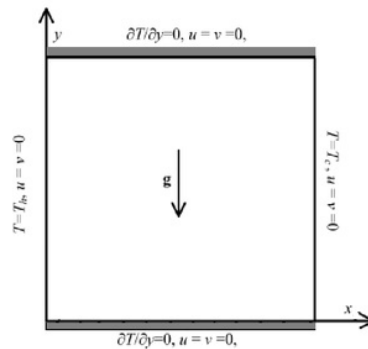


Figure1. Square enclosure and boundary conditions

15 3. Lattice Boltzmann Method

The Lattice Boltzmann originates from the lattice gas cellular automata method(Sukop and Thorne, 2006). It is based on discretization of the Boltzmann Equation with the Bhatnagar-Gross-Krook (BGK) relaxation for the collision operator. Through the Chapman Enskog expansion, the discrete Boltzmann Equation with the BGK relaxation can easily derive the Navier Stokes Equations.

Two kind of particle distribution functions are used for modeling the natural convective flows, i.e. density distribution functions f_i for calculation of density and velocity fields and energy distribution functions g_i for calculation of the temperature(Guo, Shi, and Zheng, 2002).

$$\begin{aligned} f_i(\mathbf{x}+e_i, \mathbf{x}, t+e_i, t) &= \frac{1}{f} (f_i(\mathbf{x}, t) - f_i^{eq}(\mathbf{x}, t)) + F \\ g_i(\mathbf{x}+e_i, \mathbf{x}, t+e_i, t) &= \frac{1}{g} (g_i(\mathbf{x}, t) - g_i^{eq}(\mathbf{x}, t)) \end{aligned} \quad (1)$$

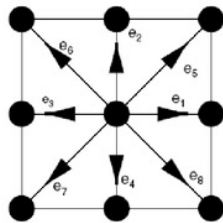


Figure 2. Discrete velocity for D2Q9

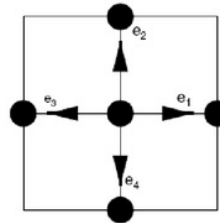


Figure 3. Discrete velocity for D2Q5 Lattice

In this paper, the D2Q9 lattice is used to represent density distribution function and D2Q5 lattice is used to represent energy distribution function. The macroscopic variables, such as density, velocity and temperature can be calculated as follow:

$$\rho = \sum_{i=0}^8 f_i ; \quad u_j = \sum_{i=0}^8 f_i e_{ij} ; \quad T_j = \sum_{i=0}^4 g_i e_{ij} \quad (2)$$

The discretized equilibrium distribution functions are given by:

$$f_i^{eq} = \frac{1}{n_i} \left(1 + 3e_i \cdot u + 4.5(e_i \cdot u)^2 - 1.5|u|^2 \right) \quad (3)$$

$$g_i^{eq} = T \cdot T_i (1 + 3e_i \cdot u_i)$$

The Boussinesq approximation is used to calculate the buoyancy force (\mathbf{F})

$$F_i = 3 \cdot n_i (e_{xi} g_{xi} + e_{yi} g_{yi}) \left(\frac{T}{T_h} - \frac{T_0}{T_c} \right) \quad (4)$$

where \mathbf{g} is gravitational acceleration, n_i is weight function and T_0 is averaged temperature.

The no slip boundary condition at the walls is implemented by using bounce back rule. For the top and bottom boundary nodes, the energy distribution functions are set equal to value g_i at nearest nodes, which are perpendicular to the wall. The Dirichlet boundary conditions of the left and right sides were implemented as follow:

```
for ( j=0; j<nj; j++){
    i0 = I2D(ni,0,j);
    i1 = I2D(ni,ni-1,j);
    T3[i1]=T_cold-T0[i1]-T1[i1]-T2[i1]-T4[i1];
    T1[i0]=T_hot-T0[i0]-T2[i0]-T3[i0]-T4[i0];}
```

3. GPU CUDA

The GPU was originally built for graphics rendering purposes, but now the GPU can be applied to accelerate numerical computation processing. In GPU CUDA programming, tasks of an application are grouped into instruction sets and passes onto the GPU such that each thread core works on different data but executes the same instruction (Thibault, 2009). In LBM methods, each node is addressed into a block thread and takes into account the value of the offset input. Basically the algorithm of LBM consists of two steps, namely, streaming and colliding. The parallel streaming step is implemented as follows:

```
3
i = blockIdx.x*TILE_I + threadIdx.x;
j = blockIdx.y*TILE_J + threadIdx.y;
i2d = i + j*pitch/sizeof(float);
f1_data[i2d] = tex2D(f1_tex, (float) (i-1), (float) j);
f2_data[i2d] = tex2D(f2_tex, (float) i, (float) (j-1));
f3_data[i2d] = tex2D(f3_tex, (float) (i+1), (float) j);
f4_data[i2d] = tex2D(f4_tex, (float) i, (float) (j+1));
...
```

The implementation of the parallel of the colliding step is:

```
3
i = blockIdx.x*TILE_I + threadIdx.x;
j = blockIdx.y*TILE_J + threadIdx.y;
i2d = i + j*pitch/sizeof(float);
if (!solid_data[i2d]==0) {

    // Read all f's and store in registers
    f0now = f0_data[i2d];
    f1now = f1_data[i2d];
    f2now = f2_data[i2d];
    ...

    // Macroscopic flow props:
    ro = f0now + f1now + f2now + f3now + f4now + f5now + f6now + f7now + f8now;
    vx = (f1now - f3now + f5now - f6now - f7now + f8now)/ro;
    vy = (f2now - f4now + f5now + f6now - f7now - f8now)/ro;

    float Ti2d = T0_data[i2d] + T1_data[i2d] + T2_data[i2d] + T3_data[i2d] + T4_data[i2d];
    T_data[i2d]= Ti2d;

    // Calculate equilibrium f's
```

```

v_sq_term = 1.5f*(vx*vx + vy*vy);
f0eq = ro * faceq1 * (1.f - v_sq_term);
f1eq = ro * faceq2 * (1.f + 3.f*vx + 4.5f*vx*vx - v_sq_term);
f2eq = ro * faceq2 * (1.f + 3.f*vy + 4.5f*vy*vy - v_sq_term);
...
force0 = 3.0*ro*faceq1*(Ti2d-T_mid)*( 0.0*gr)/(T_hot-T_cold);
force1 = 3.0*ro*faceq2*(Ti2d-T_mid)*( 0.0*gr)/(T_hot-T_cold);;
force2 = 3.0*ro*faceq2*(Ti2d-T_mid)*( 1.0*gr)/(T_hot-T_cold);;
...
// Do collisions
f0_data[i2d] = rtau1 * f0now + rtau * f0eq + force0;
f1_data[i2d] = rtau1 * f1now + rtau * f1eq + force1;
f2_data[i2d] = rtau1 * f2now + rtau * f2eq + force2;
...
T0eq = Ti2d * Tfaceq1 * (1.f +3.0*( 0.0*vx+0.0*vy));
T1eq = Ti2d * Tfaceq2 * (1.f +3.0*( 1.0*vx+0.0*vy));
....
// Simulate collisions by "relaxing" toward the local equilibrium
T0_data[i2d] = rtauT1 * T0_data[i2d] + rtauT * T0eq;
T1_data[i2d] = rtauT1 * T1_data[i2d] + rtauT * T1eq;
...

```

5. Results and Discussions

The LBM source codes are based on Graham Pullan's LBM codes(Pullan, 2012). The original codes were developed for hydrodynamical simulations. Modifications were done by adding energy distribution functions and incorporating handling of the temperature boundary conditions. The implementation of numerical calculations was done by using the Visual C++ on Intel I5 processors. The Rayleigh (Ra) numbers were varied from 1×10^3 to 1×10^6 and constant Prandtl number $Pr = 0.717$. We used 3 types of grids for all the ranges of Ra numbers. The grids consisted of 320×320 , 640×640 , 960×960 nodes.

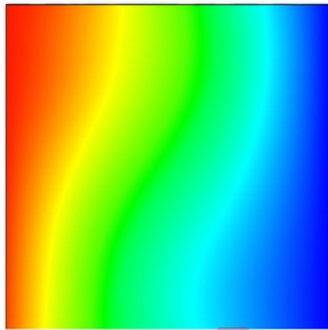


Fig. 4a. Isotherm $Ra=10^3$

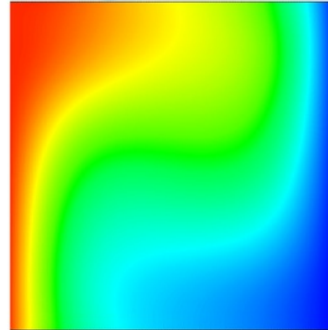


Fig. 4b. Isotherm $Ra=10^4$

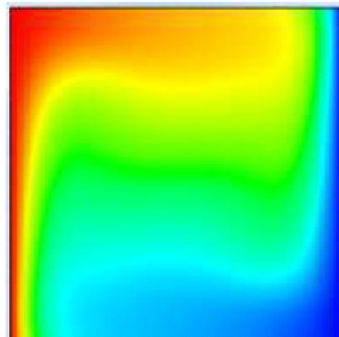


Fig. 4c. Isotherm $Ra=10^5$

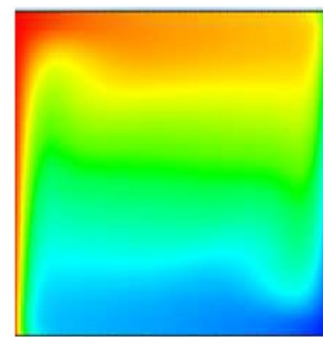


Fig. 4d. Isotherm $Ra=10^6$

The calculation results of the isotherms at steady states for different Ra are shown in Fig. 4a-4d. At low Ra number (10^3), the fluid motion driven by the buoyancy force is slow, leading to the strong diffusion. Consequently, the corresponding isotherms exhibit rather slight difference as compared to those of pure heat conduction between the left and right sides. As the Ra number is increased, the buoyancy force accelerates the circulation of fluid flow and the heat transfer of the natural convection, which is described by Nusselt Number (Nu), is significantly enhanced (figure 5).

Figure 5 shows the comparison of the present Nusselt number calculation with the literature results (Dixit and Babu, 2006; de Vahl Davis, 1983). The experiment data was taken from the empirical formulation of Martynenko and Khramtsov (Martynenko and Khramtsov, 2005). The present results agree well with other numerical results, but slightly different when compared with experiment results for high Ra number.

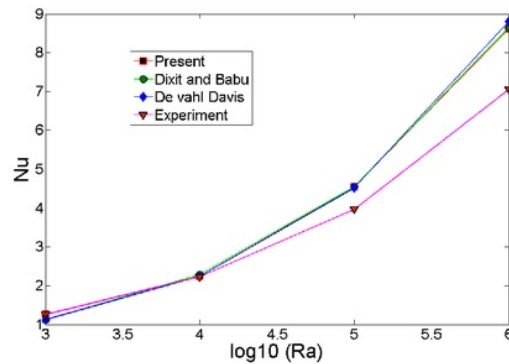


Fig.5. Comparison of the Nusselt Number

Grid independence of the present numerical results have been established for two different Rayleigh numbers, i.e. 10^4 and 10^5 . Table 1 gives the Nusselt numbers on three grids. The maximum change of the Nusselt number is 3.2 % for $Ra=10^4$ and 0.7 % for $Ra=10^5$. Those calculations show that present numerical results are grid independent.

Table 1. Grid independence study

Mesh	320×320	640×640	960×960
Nu ($Ra=10^4$)	2.241	2.244	2.158
Nu ($Ra=10^5$)	4.535	4.503	4.533

The GPU performances are compared with CPU performance $Ra=10^5$ on the three types of grids. For obtaining the steady state, all simulations were run for 30000 time steps. The GPU code performs 1.67, 6.83 and 7.53 times faster than CPU as shown in Table 2. The accelerated processing is more apparent with larger grid sizes as the intensity of arithmetic calculation and data communications in the memory of the GPU also increase, therefore aiding the calculation more and thus reducing the processing time more significantly.

Table 2. Comparison of performance for $Ra=10^5$

Grids	320×320	640×640	960×960
CPU Execution Time (s)	836.80	5129.49	11005.98
GPU Execution Time (s)	500.32	751.03	1461.86
Speed up	1.67	6.83	7.53

6. Conclusions

We have presented the Lattice Boltzmann method for solving natural convection heat transfer on GPU Nvidia GeForce 9600 card. It has been shown that the numerical results of the LBM have a good agreement with the numerical and experimental results reported in the previous studies. The numerical computation on GPU is 1.67 – 7.53 faster than on CPU and the increase in processing speed is dependant the size of the grids.

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